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# Quantum Theory and Global Optimisation

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This talk is not about Quantum Algorithms!

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Theme I:

Additivity Problems in QIT

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# Algebraic Quantum Information Theory

- Quantum Information Theory (QIT) is a two-legged science:
  - Information Theory
  - Quantum Mechanics
- “My” kind of QIT = Algebraic QIT = QIT with a third leg: tools
  - Matrix Theory: matrix inequalities, eigenvalues, singular values,...
  - Convexity Theory: convex optimisation, duality theory, convex hulls,...
- Focus on problems where these tools are essential.

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# Open Problems

- QIT is about bi- or even multipartite states
- These are essentially higher-order tensors
- Life would be easy if a higher-order SVD existed
- It does not, so we get a lot of difficult problems in QIT
  - Separability problem
  - Proving additivity of various basic quantities
- There is hope we can tackle these problems using “ordinary” linear algebra

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# Entanglement Measures

- How to quantify entanglement of a state?
- Pure states: only one reasonable measure
  - $E(|\psi\rangle\langle\psi|) = S(\text{Tr}_A |\psi\rangle\langle\psi|)$ ,  $S$  von Neumann entropy
- Mixed states: whole zoo of measures, use what you need
  - Entanglement Cost  $E_C$
  - Entanglement of Distillation  $E_D$
  - Entanglement of Formation (EoF,  $E_F$ )
  - Relative Entropy of Entanglement  $E_R$
  - Squashed Entanglement  $E_{sq}$

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# Entanglement Cost

- Calculating the entanglement cost  $E_C$  is one of the Big Open Problems of QIT.
- $E_C$  defined in an operational way, nearly impossible to calculate
- Hayden, Horodecki and Terhal:  $E_C$  is equal to the *regularisation* of  $E_F$ :

$$E_C(\rho) = \lim_{n \rightarrow \infty} E_F(\rho^{\otimes n})/n.$$

- $E_C$  would be equal to  $E_F$  if  $E_F$  were **additive**.

$$E_F(\rho_1 \otimes \rho_2) =? E_F(\rho_1) + E_F(\rho_2)$$

- This is Additivity Problem #1.
- One only needs to prove **superadditivity**,  $E_F(\rho_1 \otimes \rho_2) \geq? E_F(\rho_1) + E_F(\rho_2)$ , because **subadditivity**,  $E_F(\rho_1 \otimes \rho_2) \leq E_F(\rho_1) + E_F(\rho_2)$ , is trivial to prove.

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# Strong Superadditivity of EoF

- Vollbrecht and Werner conjectured a stronger property implying superadditivity:

**Strong Superadditivity:**  $E_F(\rho) \geq? E_F(\rho_I) + E_F(\rho_{II})$

Here  $\rho$  is a general state over a duplicated Hilbert space and  $\rho_I$  and  $\rho_{II}$  are its reductions to the different copies of that space.

- This is Additivity Problem #2
- And now for something completely different...

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# Classical Capacity of Quantum Channels

- Noisy communication channels modelled as completely positive trace-preserving (CPTP) maps between operator algebras.
- One of the most fundamental questions in QIT: determination of *classical capacity of a quantum channel* i.e. the capacity of quantum channels to transmit classical information.
- Much more difficult than its purely classical counterpart due to the existence of entanglement.
- To obtain an optimal quantum channel *decoder* one has to perform entangled measurements over the channel output states.

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# Additivity of Holevo Capacity

- Is entanglement also necessary to obtain an optimal *encoder*?
- Widely believed not to be the case, i.e. no benefit is expected in having entanglement between the (single-letter) states sent over the channel.
- To prove this, it is necessary to show that the single-letter classical capacity (a.k.a. the Holevo capacity  $\chi$ ) of a quantum channel is additive.
- The Holevo capacity of a channel  $\Phi$  is:

$$\chi(\Phi) = \sup_{\pi, \rho} S\left(\sum \pi_i \Phi(\rho_i)\right) - \sum \pi_i S(\Phi(\rho_i)).$$

- Is this additive:  $\chi(\Phi_1 \otimes \Phi_2) = \chi(\Phi_1) + \chi(\Phi_2)$ ?
- This is Additivity Problem #3.

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# Maximal Output Purity of Channels

- Holevo capacity looks quite complicated. It has resisted any attempt so far at proving its additivity.
- For that reason, a simpler channel property, called the minimal output entropy (MOE), has been introduced.
- In general, if an operation acts on a pure state, the output will be a mixed state.
- The *maximal output purity* (MOP),  $\nu_q$ , of an operation or channel quantifies how close to purity one can get by choosing the input state. As a measure of purity, one can use the Schatten  $q$ -norm  $\|X\|_q = (\text{Tr } X^q)^{1/q}$ :

$$\nu_q(\Omega) = \max_{\psi} \{ \|\Omega(|\psi\rangle\langle\psi|)\|_q : \|\psi\| = 1 \}.$$

- The *minimal output entropy* (MOE),  $\nu_S$ , is the entropic version of the MOP, where the von Neumann entropy  $S$  is used as a measure of purity.

$$\nu_S(\Omega) = \min_{\psi} \{ S(\Omega(|\psi\rangle\langle\psi|)) : \|\psi\| = 1 \}.$$

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# Additivity Problem #4

- Is MOE additive?
- Could shed some light on the additivity problem for the Holevo capacity.
- Additivity of MOE would follow from multiplicativity of MOP for “small”  $q$  (Amosov, Holevo and Werner):  $\nu_q(\Phi \otimes \Omega) = \nu_q(\Phi)\nu_q(\Omega)$ .
- Multiplicativity of MOP proven by Chris King for
  - entanglement breaking channels
  - unital qubit maps
  - depolarising channels
- Refuted for values of  $q > 4.79$  (Holevo and Werner) (Baka!)
- Nevertheless, it could still hold for  $q \downarrow 1$  (Hope springs eternal...)

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Theme II:

Optimisation Theory

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# Local and Global Optimisation

- There are lots of techniques for finding the maxima of a general function
- Most methods can get stuck in local maxima
- Global optimisation is about finding the biggest local maximum
- No method exists that guaranteedly finds that
- So we are very happy when we have a problem for which there is only one local optimum
- Convex problems are in that league.

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# Convex Sets and Functions

- A set  $S$  is **convex** if and only if

$$\forall p, q \in S : \overline{pq} \in S$$

where  $\overline{pq}$  is the line segment joining  $p$  and  $q$

- Whatever one can say about sets, one can say about functions
- The **Graph** of a function  $y = f(x)$  is the set of points  $\{(x, y) : y = f(x)\}$
- The **Epigraph** of a function  $y = f(x)$  is the set of points

$$\text{Epi}(f) := \{(x, y) : y \geq f(x)\}$$

- A function is convex if and only if its epigraph is a convex set
- A function  $f$  is **concave** if  $-f$  is convex

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# Duality Theory

- A dual view of convexity: a convex set  $S$  equals the intersection of all halfplanes containing  $S$
- A function  $f$  is convex if and only if

$$\forall \{(a_i, x_i)\} : f\left(\sum_i a_i x_i\right) \leq \sum_i a_i f(x_i)$$

- Dually,  $f$  is convex if and only if  $f$  equals the pointwise supremum of all affine functions majorised by  $f$

$$\forall x : f(x) = \sup_{a,b} \{a^T x + b : (\forall y : a^T y + b \leq f(y))\}$$

- For every  $a$ , one can calculate  $\hat{b}(a)$ , the largest  $b$  satisfying the condition.

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# Duality in Optimisation

- Let's tackle the following constrained optimisation:
- Minimise a convex function  $f(x)$  over the interval  $x \leq x_u$ .

$$\hat{f} = \min_x \{f(x) : x \leq x_u\}.$$

- We can easily prove

$$\min_x \{f(x) : x \leq x_u\} \geq \max_a \{ax_u + \hat{b}(a) : a \leq 0\}.$$

(In fact, equality holds!)

- The maximisation over  $a$  is a convex problem, called the **dual** problem.
- The minimisation over  $x$  is called the **primal** problem.

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# Certificates of Convergence

- By picking specific  $x \leq x_u$  (so-called **feasible**  $x$ ), you get upper bounds on  $\hat{f}$ .
- By picking specific tangents  $a \leq 0$ , you get lower bounds on  $\hat{f}$ .
- Thus if you solve the primal problem together with the dual one, you can bracket the solution within an upper and a lower bound.
- The difference between the bounds is an upper bound on how far you are from the real solution.
- You, therefore, get a **certificate of convergence**.
- This works for general convex optimisation problems.
- No other optimisation problem has this feature.

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# Duality and Convex Hulls

- The **Convex Hull**  $\text{Conv}(S)$  of a set  $S$  is the *union* of all line segments  $\overline{pq}$  where  $p, q \in S$
- Dually, the convex hull of  $S$  is the *intersection* of all halfplanes containing  $S$ .
- The convex closure  $\hat{f}$  of a function  $f$  is defined by

$$\text{Epi}(\hat{f}) = \text{Conv}(\text{Epi}(f))$$

- Functionally, the convex closure of  $f$  can be calculated by taking convex combinations

$$\hat{f}(x) = \min_{\{(a_i, x_i)\}} \left\{ \sum_i a_i f(x_i) : \sum_i a_i x_i = x \right\}$$

- And, dually, the convex closure of  $f$  is the pointwise supremum of all affine functions majorised by  $f$

$$\hat{f}(x) = \sup_{a, b} \{ a^T x + b : (\forall y : a^T y + b \leq f(y)) \}$$

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Allegro:

Convex Optimisation and Additivity

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# Entanglement of Formation Defined

- Any state  $\rho$  can be realised by an **ensemble** of pure states
- An ensemble is specified by a set of pairs  $\{(p_i, \psi_i)\}_{i=1}^N$ 
  - of  $N$  state vectors  $\psi_i$  and statistical weights  $p_i$
  - with  $p_i \geq 0$  and  $\sum_i p_i = 1$
- The entanglement of formation (EoF) of a bipartite state  $\rho$  (over the bi-partite Hilbert space  $\mathcal{H}_A \otimes \mathcal{H}_B$ ), is

$$E_F(\rho) = \min_{\{(p_i, \psi_i)\}} \left\{ \sum_i p_i S(\text{Tr}_A \Psi_i) : \sum_i p_i \Psi_i = \rho \right\}.$$

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# EoF is a Convex Closure

- **Observation 1:** The definition of the EoF really means that the EoF is the **convex closure** of the pure state entanglement function

$$E(\Psi) = S(\text{Tr}_A \Psi)$$

- Consider bounded functions  $f$  whose domain is the set of states  $\mathcal{S}(\mathcal{H})$
- We can apply real convex analysis because  $\mathcal{B}(\mathcal{H})$  with  $\langle A, B \rangle = \text{Tr}[AB]$  is a real vector space
- **Definition 1:** The convex closure of  $f$  is

$$\hat{f}(\rho) = \min_{\{(p_i, \rho_i)\}} \left\{ \sum_i p_i f(\rho_i) : \sum_i p_i \rho_i = \rho \right\}$$

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# EoF is a Convex Closure

- **Observation 2:** The second, dual, formulation of the convex closure looks easier
- **Definition 2:** The convex closure of a function  $f$  is the pointwise supremum of all affine functions on  $\mathcal{S}(\mathcal{H})$  majorised by  $f$

$$\hat{f}(\rho) = \sup_{X \in \mathcal{B}(\mathcal{H})} \{ \text{Tr}[\rho X] : (\forall \sigma \in \mathcal{S}(\mathcal{H}) : \text{Tr}[\sigma X] \leq f(\sigma)) \}$$

- Moreover, this definition can be pulled apart into two identical parts, based on the concept of the **conjugate function**

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# Conjugate function

- Define the **conjugate function**  $f^*$ :

$$f^*(X) = \max_{\rho \in \mathcal{S}(\mathcal{H})} \text{Tr}[\rho X] - f(\rho)$$

- If  $f$  is continuous this is called the **Legendre transform** of  $f$ .
- The conjugate is convex in  $X$ : pointwise maximum of affine functions
- The conjugate and convex closure determine each other completely

$$f \xrightarrow{*} f^* \xleftarrow{*} \hat{f}$$

- The convex closure of  $f$  is the conjugate of the conjugate of  $f$ :  $\hat{f} = f^{**}$
- The conjugate of the convex closure of  $f$  is the conjugate of  $f$ :  $\hat{f}^* = f^*$

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# Additivity of EoF

**Theorem 1 (Audenaert and Braunstein)** *With  $g$  defined by*

$$g(A) = \max_{\psi \in \mathcal{H}} \text{Tr}[\Psi \log(A)] - E(\Psi)$$

*strong superadditivity of the EoF,*

$$E_F(\rho) \geq? E_F(\rho_I) + E_F(\rho_{II}),$$

*is equivalent to subadditivity of  $g$ ,*

$$g(A_1 \otimes A_2) \leq? g(A_1) + g(A_2).$$

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# Relation to MOP

- The above Theorem reduces the additivity problem for the EoF, defined as a minimisation over *ensembles*, to an equivalent problem for the conjugate function, defined as a maximisation over *pure states*.
- Using the Lie-Trotter relation, entropic quantities can be converted to power-law quantities.
- Using some involved mathematics, Audenaert and Braunstein also proved:  
**Theorem 2** *If  $\nu_q$  is multiplicative for  $q \downarrow 1$  and for all completely positive maps, then the entanglement of formation is strongly superadditive.*
- So two additivity problems are closely related!

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# Equivalence of Additivity Problems

- One month later...
- Peter Shor proved equivalence of all four additivity problems!
  - additivity of EoF,
  - strong superadditivity of EoF,
  - additivity of classical capacity of a quantum channel,
  - additivity of the MOE:  $\nu_S(\Phi \otimes \Omega) = \nu_S(\Phi) + \nu_S(\Omega)$ .
- The stakes for proving multiplicativity of MOP have raised.
- So what's next?

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Theme III:

The Quantum de Finetti Theorem

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# Unknown Probabilities

- Probability Estimation is a procedure to estimate an unknown probability from the results of repeated trials in identical circumstances.
- In Bayesian view of probability theory, a probability is a “measure of credible belief, reflecting one’s state of knowledge”.
- In this view, “unknown probability” is an oxymoron.
- Bruno de Finetti (early 1930’s) tried to eliminate this offending concept.
- He focused on the equivalence of repeated trials:
  - indistinguishability of the different trials w.r.t. predictions
  - a probability assignment for multiple trials should be symmetric under permutation of the trials

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# Exchangeable Probabilities

- Principle of Equivalence of Repeated Trials:
  - If an experimenter judges a collection of  $N$  trials to have a probability  $P^{(N)}$ , he will judge any permutation of the trials to have that same probability.
  - This will be true for every  $N$
  - Consistency condition:  $P^{(N)}$  must be derivable from  $P^{(N+1)}$
- Sequence of probabilities  $(P^{(N)})_N$  obeying this condition are called **exchangeable**.

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# de Finetti's Theorem

- A sequence of probabilities  $(P^{(N)})_N$  is exchangeable if and only if

$$P^{(N)} = \int d\mu(P) P^{\times N},$$

where  $d\mu(P)$  is a positive measure over (single-trial) probabilities.

- “Probability distribution over probabilities” replaces “unknown probability”
- Experimenter can act as if...
  - there is an objective (single-trial) probability assignment,  $P$ ,
  - yielding  $P^{\times N}$  for  $N$  repeated trials,
  - his uncertainty about  $P$  is expressed by  $\int d\mu(P)$ .

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# Exchangeable Quantum States

- In Quantum Theory, the analogon of a probability is the *quantum state*
- In quantum tomography, one tries to measure an unknown quantum state using repeated trials on identically prepared particles.
- In the information-based interpretation of quantum mechanics, a state represents the state of knowledge of an observer.
- Again, in that interpretation, “unknown quantum state” is an oxymoron.
- Sequence of **exchangeable quantum states**  $(\rho^{(N)})_N$ :
  - State  $\rho^{(N)}$  defined over  $N$ -fold copy of Hilbert space  $\mathcal{H}^{\otimes N}$ ;
  - Every  $\rho^{(N)}$  is symmetric under permutation of copies of  $\mathcal{H}$
  - Consistency:  $\rho^{(N)} = \text{Tr}_1 \rho^{(N+1)}$

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# The Quantum de Finetti Theorem

- Theorem (Hudson and Moody 1976; Størmer 1969):

A sequence of states  $(\rho^{(N)})_N$  is exchangeable if and only if

$$\rho^{(N)} = \int d\mu(\rho) \rho^{\otimes N},$$

where  $d\mu(\rho)$  is a positive measure over the state space of  $\mathcal{H}$ .

- Concept of “unknown state” replaced by “probability distribution over states”
- Tomographer can act as if...
  - there is an “observer-in-the-box” preparing systems in the same state  $\rho$ ,
  - yielding  $\rho^{\otimes N}$  for  $N$  repeated trials,
  - his uncertainty about  $\rho$  is expressed by  $\int d\mu(\rho)$ .

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*Allegretto:*

*There, and Back Again*

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# The Plan

- Optimisation theory can help us in solving questions in quantum theory.
- It has done so before.
- It can't help us very much with the MOP:
- Not a convex problem: it is a *maximisation* of a convex function over a convex set
- All we get from convexity theory is that the maximum will be obtained in an extreme point
- Now here is where Quantum de Finetti comes in!
- So, in return, quantum theory can offer solutions to questions in optimisation theory!

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# Maximal Output Purity (again!)

- Consider the maximal output purity  $\nu_q$  for integer  $q$ ; then

$$\nu_q^q(\Phi) = \max_{\rho \in \mathcal{S}(\mathcal{H})} \text{Tr}[(\Phi(\rho))^q],$$

- Note that  $\text{Tr}[(\Phi(\rho))^q] = \text{Tr}[\Phi(\rho)\Phi(\rho) \dots \Phi(\rho)]$ , with  $q$  factors.
- We can write  $\text{Tr}[(\Phi(\rho))^q] = \text{Tr}[A\rho^{\otimes q}]$ , with

$$A_{(i),(j)} = \text{Tr}[\Phi_{i_1,j_1} \dots \Phi_{i_q,j_q}],$$

where  $\Phi$  is the Choi matrix of the map  $\Phi$ .

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# Symmetry

- Consider permutations of  $n$  copies of  $\mathcal{H}$ ,  $\pi \in S_n$ .
- If, for every permutation  $\pi \in S_n$ , a matrix  $A$  over  $\mathcal{H}^{\otimes n}$  obeys

$$A_{(i_1, \dots, i_q), (j_1, \dots, j_q)} = A_{(i_{\pi(1)}, \dots, i_{\pi(q)}), (j_{\pi(1)}, \dots, j_{\pi(q)})},$$

then the matrix  $A$  is *symmetric*.

- Let  $P_\pi$  permute the indices according to  $\pi$ , i.e.  $(P_\pi x)_{(i)} = x_{\pi(i)}$ .
- Thus  $A$  is symmetric if and only if  $\forall \pi \in S_n, P_\pi^\dagger A P_\pi = A$ .
- The linear map  $P_n$  that projects all operators to the symmetric subspace is

$$P_n(A) = \frac{1}{n!} \sum_{\pi \in S_n} P_\pi^\dagger A P_\pi.$$

We call  $P_n(A)$  the *symmetric part* of  $A$ .

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# Yet Another Main Theorem

The symmetric part of  $A$  is:

$$P_n(A) = \frac{1}{n!} \sum_{\pi \in S_n} P_\pi^\dagger A P_\pi.$$

**Theorem 3** *For any  $q, n \in \mathbb{N}$ , and for any operator  $A$  over  $\mathcal{H}^{\otimes q}$  with Hermitian symmetric part, the sequence  $(\mu_n(A))_n$ , with*

$$\mu_n(A) := \lambda_{\max}(P_{q+n}(A \otimes \mathbf{I}^{\otimes n})),$$

*is non-increasing and converges to*

$$\lim_{n \rightarrow \infty} \mu_n(A) = \max_{\rho} \text{Tr}[A\rho^{\otimes q}].$$

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# Proof of Theorem (I)

- We want to maximise  $\text{Tr}[A\rho^{\otimes q}]$  over all states  $\rho$ .
- First turn this into the more general optimisation problem

$$\max_{d\mu(\rho) \geq 0} \left\{ \text{Tr} \left[ A \int \rho^{\otimes q} d\mu(\rho) \right] : \int d\mu(\rho) = 1 \right\}$$

- We can replace this by a maximisation of  $\text{Tr}[A\rho^{(q)}]$  over all  $\rho^{(q)}$  that are  $q$ -th element in some exchangeable sequence of states:
  - Any state of the form  $\rho^{\otimes q}$  obviously forms part of an exchangeable sequence.
  - Conversely, by the QdF theorem, the  $q$ -th element in any exchangeable sequence must be of the form  $\int \rho^{\otimes q} d\mu(\rho)$  for some positive measure  $d\mu(\rho)$ .

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## Proof of Theorem (II)

- From the definition of exchangeable sequence, we infer that  $\rho^{(q)}$  must be the partial trace of a symmetric state  $\rho^{(\infty)}$  over  $\mathcal{H}^{\otimes \infty}$ , where all but  $q$  copies of  $\mathcal{H}$  have been traced out.
- As a result, we can express the states  $\int \rho^{\otimes q} d\mu(\rho)$  in

$$\max_{\rho} \text{Tr}[A\rho^{\otimes q}] = \max_{d\mu(\rho) \geq 0} \left\{ \text{Tr}\left[A \int \rho^{\otimes q} d\mu(\rho)\right] : \int d\mu(\rho) = 1 \right\}$$

as a partial trace of symmetric states over  $\mathcal{H}^{\otimes \infty}$ :

$$\begin{aligned} \max_{\rho} \text{Tr}[A\rho^{\otimes q}] &= \lim_{n \rightarrow \infty} \max_{\rho} \left\{ \text{Tr}[A \text{Tr}_n[\rho]] : \rho \text{ symmetric over } \mathcal{H}^{\otimes (q+n)} \right\} \\ &= \lim_{n \rightarrow \infty} \max_{\rho} \left\{ \text{Tr}[(A \otimes \mathbf{I}^{\otimes n}) \rho] : \rho \text{ symmetric over } \mathcal{H}^{\otimes (q+n)} \right\}. \end{aligned}$$

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# Proof of Theorem (III)

- So we have

$$\max_{\rho} \text{Tr}[A\rho^{\otimes q}] = \lim_{n \rightarrow \infty} \max_{\rho} \{ \text{Tr}[(A \otimes \mathbf{I}^{\otimes n}) \rho] : \rho \text{ symmetric} \}.$$

- This maximisation over *symmetric* states can be replaced by a maximisation over *all* states, provided they are projected first onto the symmetric subspace:

$$\max_{\rho} \text{Tr}[A\rho^{\otimes q}] = \lim_{n \rightarrow \infty} \max_{\rho} \text{Tr}[(A \otimes \mathbf{I}^{\otimes n}) \mathbf{P}_{q+n}(\rho)].$$

- The projection can equally well be applied to the factor  $A \otimes \mathbf{I}^{\otimes n}$ :

$$\begin{aligned} \max_{\rho} \text{Tr}[A\rho^{\otimes q}] &= \lim_{n \rightarrow \infty} \max_{\rho} \text{Tr}[\mathbf{P}_{q+n}(A \otimes \mathbf{I}^{\otimes n}) \rho] \\ &= \lim_{n \rightarrow \infty} \lambda_{\max}(\mathbf{P}_{q+n}(A \otimes \mathbf{I}^{\otimes n})). \end{aligned}$$

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# Proof of Theorem (IV)

- We only have to show that  $\mu_n(A)$  is non-increasing with  $n$ .
- To do so, we use the convexity of  $\lambda_{\max}$ :

$$\begin{aligned}\mu_{n+1}(A) &= \lambda_{\max}(\mathbf{P}_{q+n+1}(A \otimes \mathbf{I}^{\otimes(n+1)})) \\ &= \lambda_{\max}(\mathbf{P}_{q+n+1}(\mathbf{P}_{q+n}(A \otimes \mathbf{I}^{\otimes n}) \otimes \mathbf{I})) \\ &\leq \frac{1}{(q+n+1)!} \sum_{\pi \in S_{q+n+1}} \lambda_{\max}(P_{\pi}^{\dagger}(\mathbf{P}_{q+n}(A \otimes \mathbf{I}^{\otimes n}) \otimes \mathbf{I})P_{\pi}) \\ &= \lambda_{\max}(\mathbf{P}_{q+n}(A \otimes \mathbf{I}^{\otimes n}) \otimes \mathbf{I}) \\ &= \lambda_{\max}(\mathbf{P}_{q+n}(A \otimes \mathbf{I}^{\otimes n})) \\ &= \mu_n(A).\end{aligned}$$

- This finishes the Proof.

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# So what is this?

- Algorithmic issues: with some more symmetry theory, we get a modestly efficient algorithm that calculates **guaranteed** upper bounds and has no problems at all with local maxima!
- See [quant-ph/0402076](#) for more details than you'd care for
- There might be more in store here!
- Theoretical Issues: reduces a difficult optimisation problem to an eigenvalue problem
- This might just be the simplification needed for tackling additivity
- (Watch out for forthcoming stuff on ArXiv)